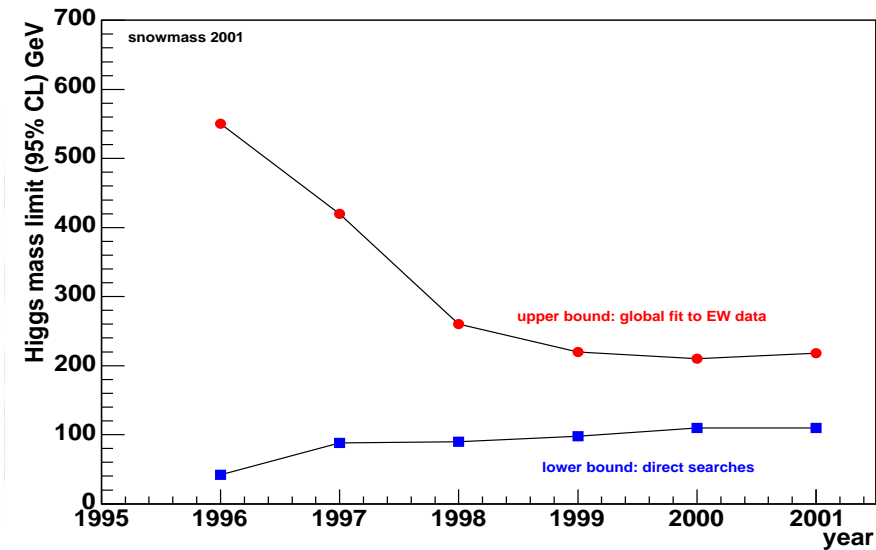
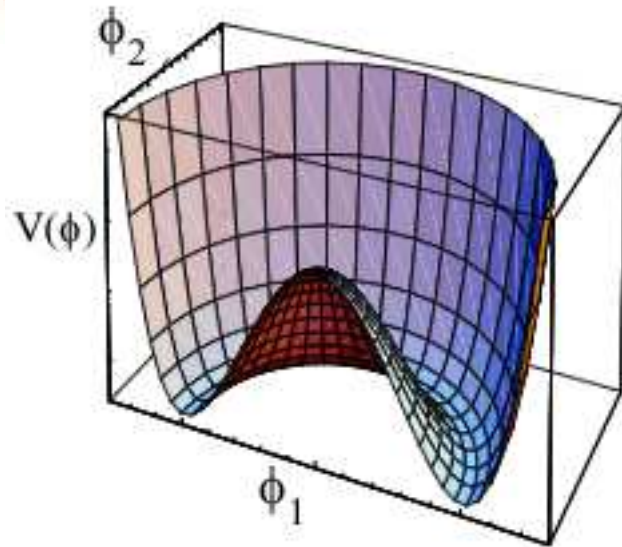


$W_L W_L$  scattering at LHC:  
(i) An approach to the Theoretical Background.  
(ii) My Performance Studies.

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University College London

- Standard Model: A very good model satisfying theorists and experimentalists.
- It explains the **Electroweak Symmetry Breaking-EWSB** by introducing the **Higgs** boson.



- However, any assumptions and any mass limits are **model dependent**.
- Alternative models to explain EWSB: SM with heavy higgs, technicolor, composite models etc etc
- Enhanced production of **longitudinal** vector boson pairs is one of the most characteristic signals of the new physics

- Describes the low energy effects of different strongly interacting models of the Symmetry Breaking Sector.
- The differences among underlying theories appear through the values of the effective chiral couplings.
- It includes operators up to order of  $s^2$  ( $E^4$ ).
- At the lowest order:

$$\mathcal{L}^{(2)} = \frac{u^2}{4} \text{Tr}\{D_\mu U D^\mu U^\dagger\} \quad (1)$$

where

$$D_\mu U = d_\mu U - W_\mu U + U B_\mu$$

$$W_\mu = -ig \frac{\sigma^\alpha W_\mu^\alpha}{2} \quad B_\mu = ig \frac{\sigma^3 B_\mu}{2} \quad U = \exp\left(\frac{i\omega^\alpha \sigma^\alpha}{u}\right)$$

where  $\sigma$  are the Pauli matrices,  $\omega$  are the three Goldstone bosons and  $u = 246$  GeV

- The next term includes the model-dependent effective couplings:

$$\mathcal{L}^{(4)} = \alpha_4 \left(\text{Tr}\{D_\mu U D^\mu U^\dagger\}\right)^2 + \alpha_5 \left(\text{Tr}\{D_\mu U D^\nu U^\dagger\}\right)^2 \quad (2)$$

- The  $\alpha_4$  and  $\alpha_5$  depend on the model but also on the renormalization scale  $\mu$ . With  $\mu = 1$  TeV we expect them to be in the range of  $[-0.01, 0.01]$
- Additional terms of the order of  $s^2$  contribute to anomalous trilinear couplings between vector bosons.

- For the  $W_L^a W_L^b \rightarrow W_L^c W_L^d$  in the weak isospin space:

$$\mathcal{M}(W_L^a W_L^b \rightarrow W_L^c W_L^d) \equiv A(s, t, u) \delta^{ab} \delta^{cd} + A(t, s, u) \delta^{ac} \delta^{bd} + A(u, t, s) \delta^{ad} \delta^{bc} \quad (3)$$

where the key amplitude  $A(s, t, u)$  is:

$$A(s, t, u) = \frac{s}{u^2} + \frac{1}{4\pi u^4} (2\alpha_5 s^2 + \alpha_4 (t^2 + u^2)) + \frac{1}{16\pi^2 u^4} \left( -\frac{t}{6} (s + 2t) \log \left( -\frac{t}{\mu^2} \right) - \frac{u}{6} (s + 2u) \log \left( -\frac{u}{\mu^2} \right) - \frac{s^2}{2} \log \left( -\frac{s}{\mu^2} \right) \right) \quad (4)$$

(In the above:  $W_L$  denotes either  $W_L^\pm$  or  $Z_L$ ,  $W_L^\pm = \frac{W_L^1 \mp i W_L^2}{\sqrt{2}}$ ,  $Z_L = W_L^3$ ,  $a, b, c, d = 1, 2, 3, 4$  and  $s, t, u$  are the Mandelstam kinematical variables.)

- Precise measurement of the  $W_L W_L \rightarrow W_L W_L$  scattering cross-section would allow the extraction of the  $\alpha_4$  and  $\alpha_5$  parameters.

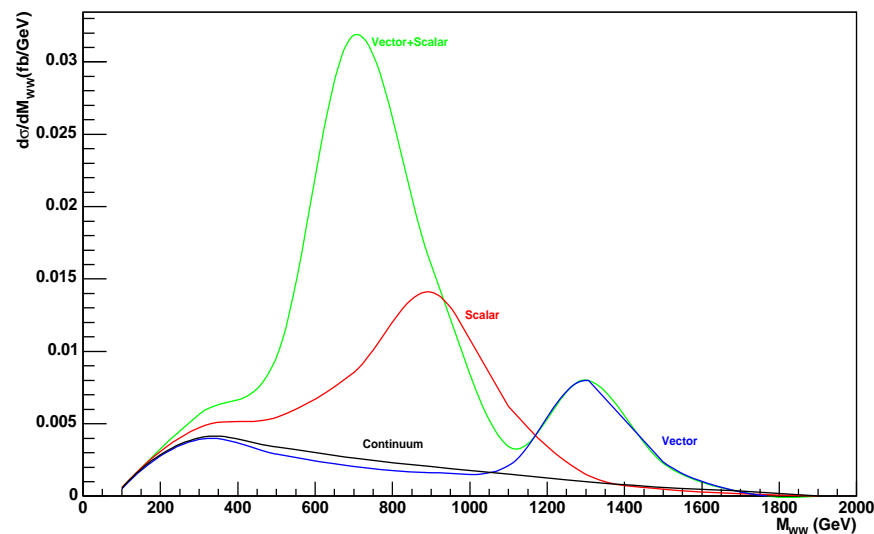
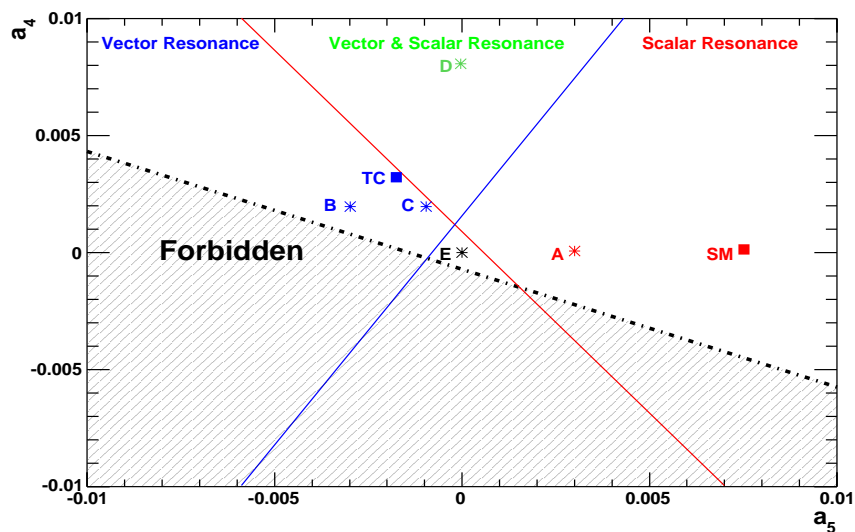
- The usual EWChL approach doesn't respect **unitarity**.
- Unitarity is restored by applying different **unitarization protocols**, for example: **Inverse Amplitude Method (Pade)**, N/D protocol etc.
- The **position** and the **nature** of the resonances **depend strongly** upon the unitarisation procedure.  
(see for example Phys.Rev.D **65** 096014 for comparison between the Pade and the N/D protocols)

- Using the Pade protocol, we obtain the following mass and the width of the resonances:

$$M_V^2 = \frac{u^2}{4(\alpha_4 - 2\alpha_5) + \frac{1}{144\pi^2}}, \quad \Gamma_V = \frac{M_V^3}{96\pi u^2} \quad (5)$$

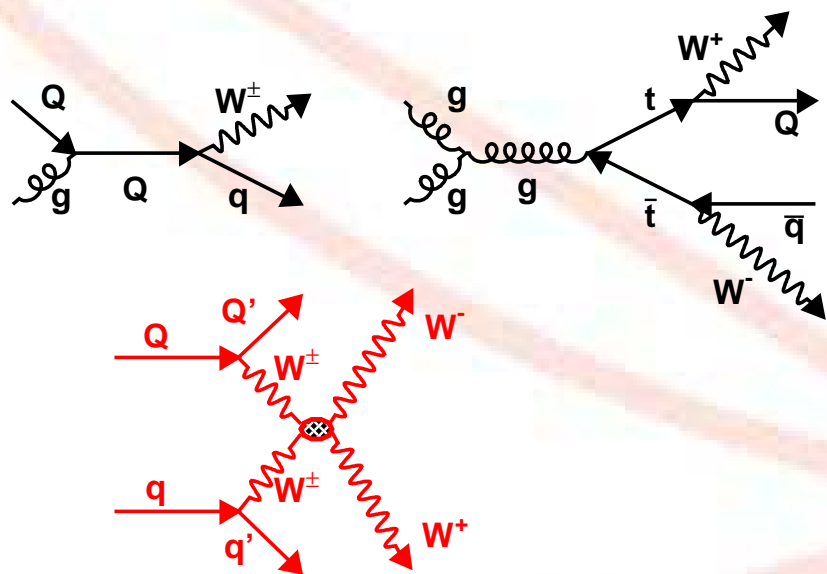
$$M_S^2 = \frac{12u^2}{16(11\alpha_5 + 7\alpha_4) + \frac{101}{48\pi^2}}, \quad \Gamma_S = \frac{M_S^3}{16\pi u^2} \quad (6)$$

- For equal masses, scalar resonances would be **6 times wider** than vector resonances.



Scenario	$\alpha_4$	$\alpha_5$	Resonance Mass (GeV)
Scalar(A)	0.0	0.003	989.8
Vector(B)	0.002	-0.003	1360.3
Scalar + Vector (D)	0.008	0.0	809.6 + 1360.3
Continuum (E)	0.0	0.0	NA

- PYTHIA has been modified to include the EWChL and to produce the resonances for different parameters according to the Pade protocol.
- Forbidden region: Where causality is violated, i.e.  $\frac{32}{3}(\alpha_5 + 2\alpha_4) + \frac{273}{864\pi^2} < 0$



- Signal:

- Private production of 600k Continuum Signal events.
- Semi-leptonic decays of the  $W$ .
- $\sigma \times BR = 3.32 fb$ .

- $t\bar{t}$  Background:

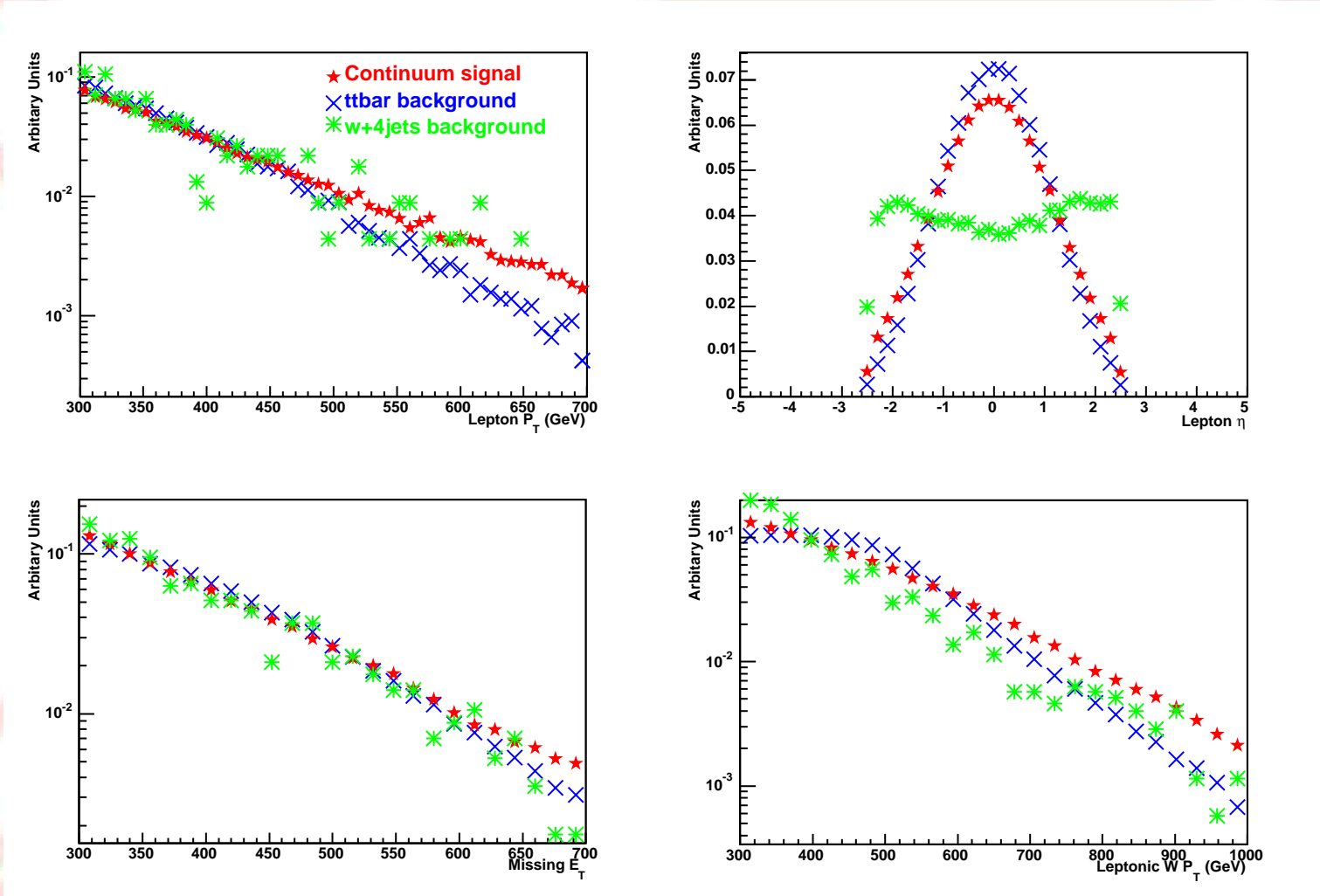
- 500k events used for Rome (T2 Sample).
- MC@NLO - Fully hadronic events were rejected and at least 1 top has  $P_T > 500 GeV$ .
- $\sigma \times BR = 4.12 fb$ .

- $w+4jets$  Background:

- 200k events used for DC2 (A7 Sample).
- AlpGen (Hard Process + Matrix Element: require 4 jets) - Pythia (Parton Shower + Hadronisation).
- Light Jets:  $P_T > 20 GeV$   $\eta_{max} = 5$   $\Delta R_{jj} = 0.4$ .
- Lepton:  $P_T > 20 GeV$   $\eta_{max} = 3$   $\Delta R_{lept-j} = 0.4$   $E_T^{miss} > 15 GeV$ .
- $\sigma \times BR = 1.20 fb$ .

The Generated Events were then Simulated using ATLFast.

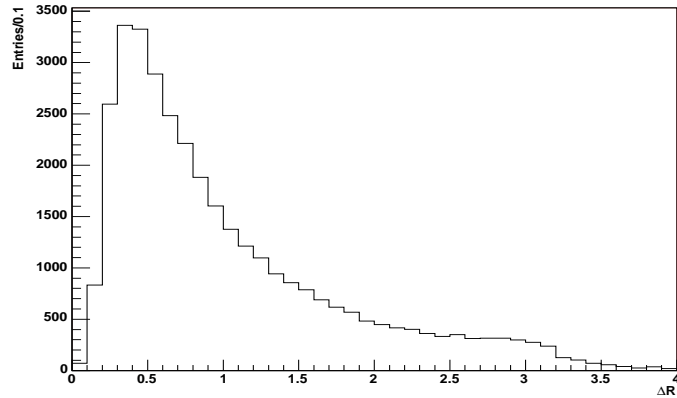
# Initial Distributions: The Leptonic sector.



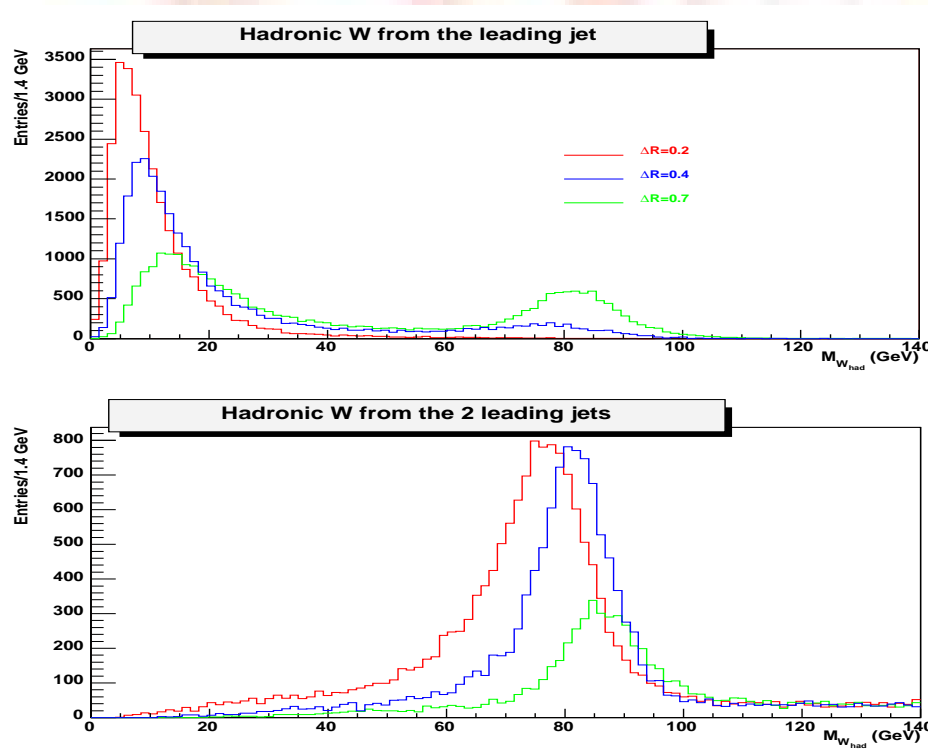
- Applied Cuts:  $P_T^{lept} > 100 GeV$ ;  $P_T^{W_{lept}} > 320 GeV$ ;  $E_T^{miss} > 100 GeV$



# Initial Distributions: The Case of the Hadronic W reconstruction.

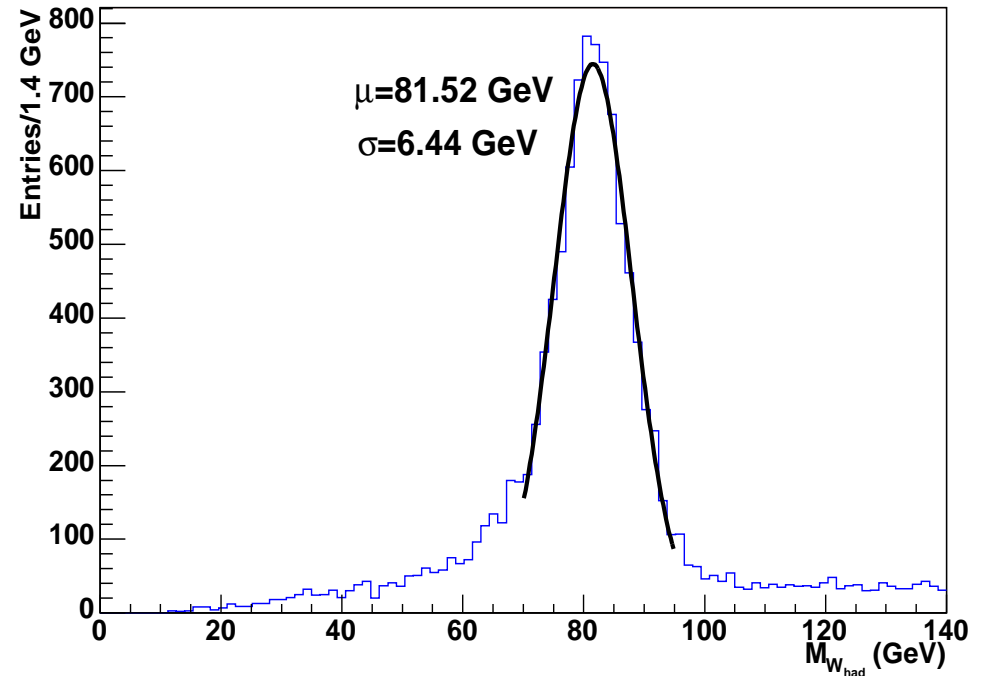
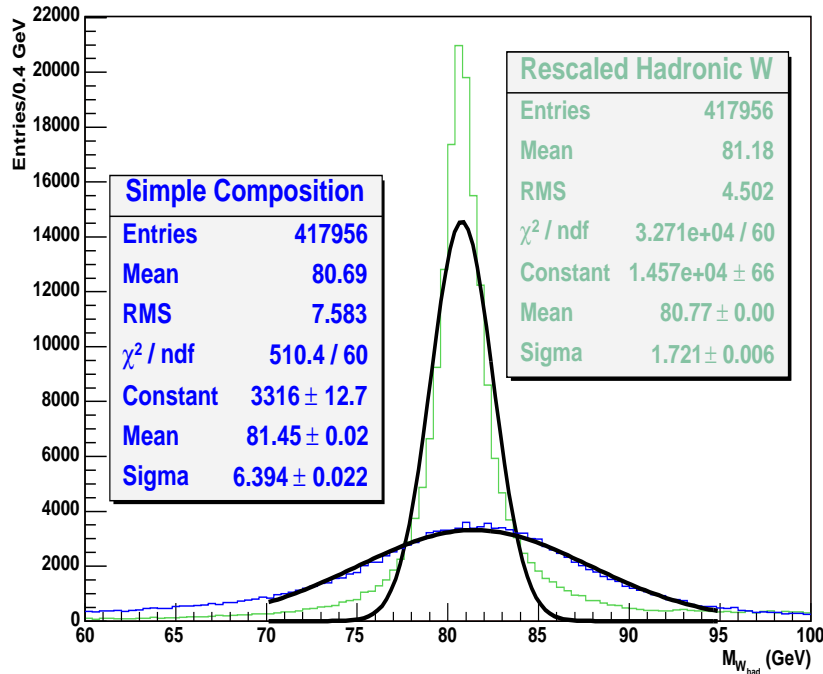


- Due to the high boost of the W, the 2 jets can be very close ( $MPV \sim 0.5$ ) and will overlap.
- The Cone Algorithm has been used for the jet finding.
- $\Delta R = 0.4$  is the optimum size for reconstructing the hadronic W from 2 jets.

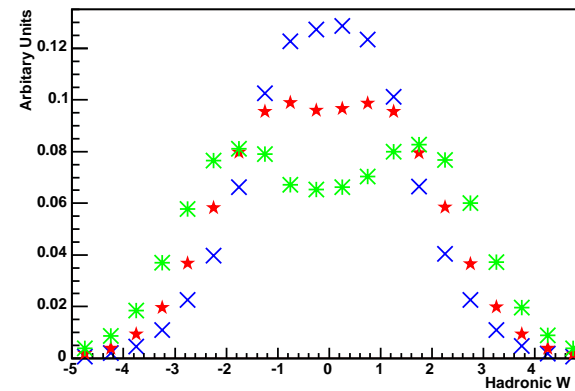
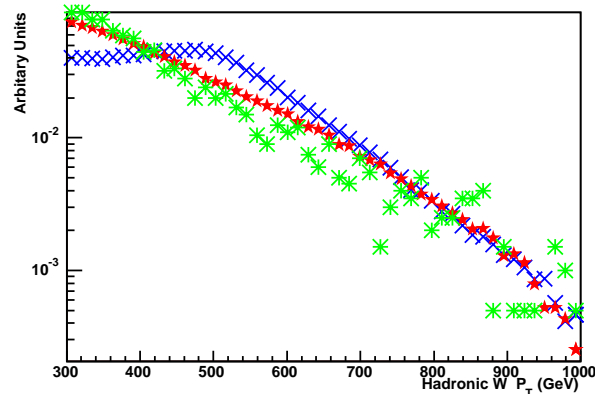
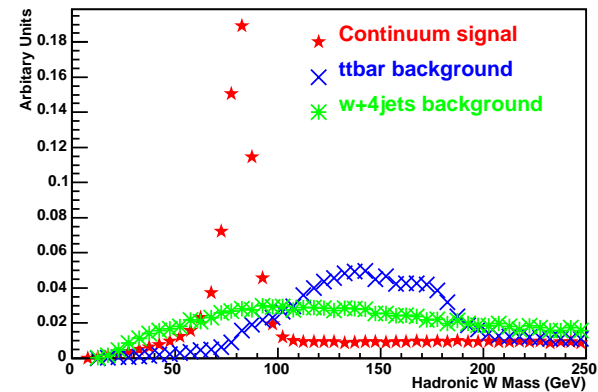
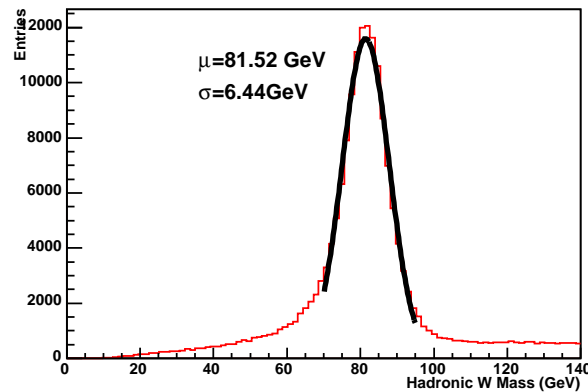


# Initial Distributions: The Case of the Hadronic W reconstruction.

Checking Simon's tool for rescaling the hadronic W (see talk on 4th November 2005).



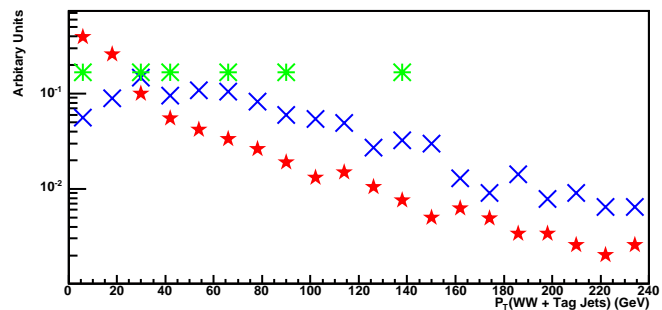
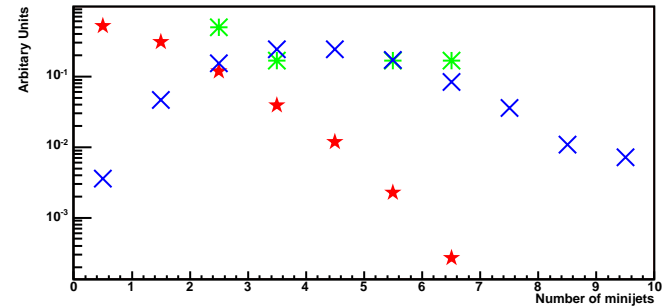
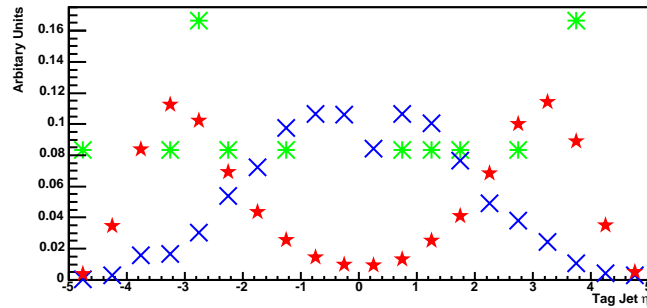
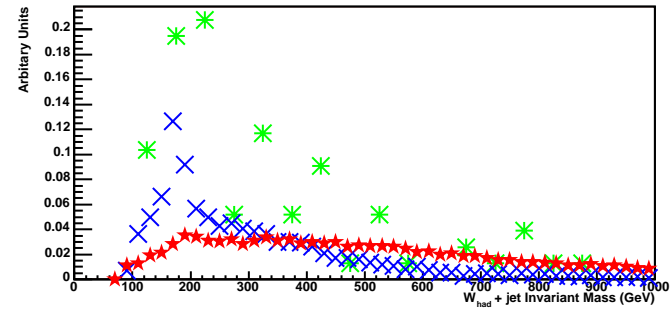
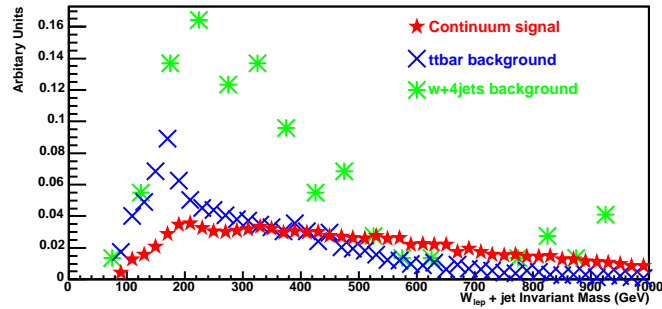
- The rescaling method gives  $\Gamma_W \ll \Gamma_W^{natural} = 2.06 \text{ GeV}!!$
- Use the simple composition of the 2 jets



- Applied Cuts:  $P_T^{W_{had}} > 320\text{GeV}$ ;  $67\text{GeV} < M_{W_{had}} < 97\text{GeV}$ ;  $\eta_{W_{had}} > 4$
- The cut on the mass of the hadronic W kills the majority of the background events.
- Some distributions appear different from those in the note!!!

# Characteristics of the Hadronic environment.

After applying the previous cuts, we investigate the features of the hadronic environment:



Cuts:

- (a-b) Top quark veto:  $130 \text{ GeV} < M_{W_{\text{jet}}} < 240 \text{ GeV}$
- (c) Tag Jets: Both in Forward/Backward regions,  $P_T^{\text{jet}} > 20 \text{ GeV}$ ;  $E^{\text{jet}} > 300 \text{ GeV}$ ;  $|\eta| > 2$
- (d) Hard  $P_T$ :  $P_T^{\text{WW+tag jets}} < 50 \text{ GeV}$
- (e) MiniJet veto: Number of minijets  $< 1$

- A theoretical introduction to the EWChL and its dynamics to describe different strongly interaction models of the EWSB has been given.
- My (very initial) analysis on the  $W_L W_L$  scattering at high masses has been presented.
- The analysis has taken into account the effects of the **multi-parton interactions (Rome tuning)** but **not pile-up** effects.
- The **Cone** Algorithm has been used (only for convenience reasons). Migrate to **KtJet** and apply the **subject** analysis.
- The signal appears to have the expected behaviour. It not always the case for the background samples (especially w+4jets).
- Problems may be due to:
  - The **analysis code**: Some improvements still to be done together the subject analysis.
  - The **background samples**: Different from what has been used in the note. Wait for the new data production (DC3).
- The exotics group has proposed to have samples of W+xJets ( $x=1,\dots,5$ ) with  $P_T^{jet} > 40\text{GeV}$  and  $P_T^{jet} > 100\text{GeV}$ . UCL (sstef) will be responsible for this production (at least for the high energy samples).